

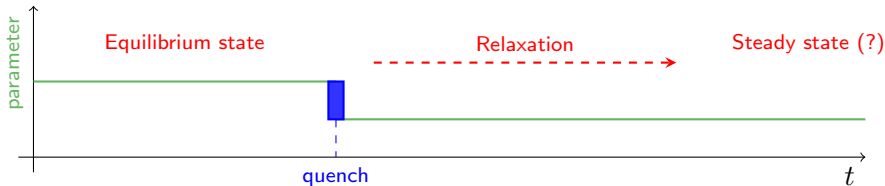


**UNIVERSITÀ
DEGLI STUDI
DI GENOVA**



Transport properties as a tool to study quench-induced dynamics in 1D systems

Definition: a **change in time** of the **parameter(s)** that govern the dynamics of an **isolated quantum system** (i.e. **under unitary time-evolution**).



Convenient way to bring a system **out-of-equilibrium** and to study its following **relaxation** toward a steady state (?).

Polkovnikov, RMP (2007)
D'Alessio et al., Adv Phys, (2014)

Definition: a **change in time** of the **parameter(s)** that govern the dynamics of an **isolated quantum system** (i.e. **under unitary time-evolution**).

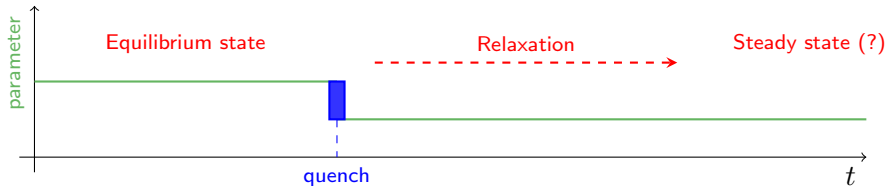


Convenient way to bring a system **out-of-equilibrium** and to study its following **relaxation** toward a steady state (?).

Polkovnikov, RMP (2007)
D'Alessio et al., Adv Phys, (2014)

- Does the steady state exist? If yes, which kind of state is it?

Definition: a **change in time** of the **parameter(s)** that govern the dynamics of an **isolated quantum system** (i.e. **under unitary time-evolution**).



Convenient way to bring a system **out-of-equilibrium** and to study its following **relaxation** toward a steady state (?).

Polkovnikov, RMP (2007)
D'Alessio et al., Adv Phys, (2014)

- Does the steady state exist? If yes, which kind of state is it?
- Which are the features of the relaxation dynamics?

Low energy physics of 1D
interacting fermion systems

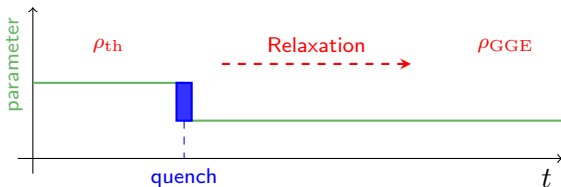
Luttinger liquid model (LL)

Integrable system:

- complete set of conserved quantities;
- a steady state is reached (in the thermodynamical limit) but it retains a **strong memory** of the initial state;
- local observables relax to **non-thermal values** described by the GGE density matrix

$$\rho_{\text{GGE}} = \frac{e^{-\sum_m \lambda_m I_m}}{Z_{\text{GGE}}}$$

Rigol et al. PRL (2007)



Low energy physics of 1D
interacting fermion systems

Luttinger liquid model (LL)

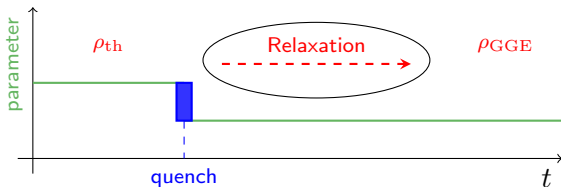
Integrable system:

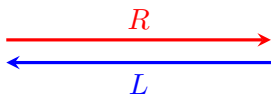
- complete set of conserved quantities;
- a steady state is reached (in the thermodynamical limit) but it retains a **strong memory** of the initial state;
- local observables relax to **non-thermal values** described by the GGE density matrix

$$\rho_{\text{GGE}} = \frac{e^{-\sum_m \lambda_m I_m}}{Z_{\text{GGE}}}$$

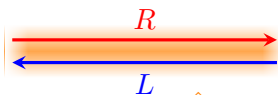
Rigol et al. PRL (2007)

Our goal: key features of
the relaxation process
towards GGE in LL



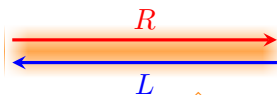


$$\hat{\mathcal{H}}_0(x) = iv_F \left(\hat{\psi}_R^\dagger(x) \partial_x \hat{\psi}_R(x) - \hat{\psi}_L^\dagger(x) \partial_x \hat{\psi}_L(x) \right)$$



The diagram shows two horizontal arrows representing modes in a Luttinger liquid. The top arrow is red and points to the right, labeled with a red 'R'. The bottom arrow is blue and points to the left, labeled with a blue 'L'. The arrows are set against a light orange gradient background.

$$\hat{\mathcal{H}}_0(x) = iv_F \left(\hat{\psi}_R^\dagger(x) \partial_x \hat{\psi}_R(x) - \hat{\psi}_L^\dagger(x) \partial_x \hat{\psi}_L(x) \right)$$
$$\hat{\mathcal{H}}_{int}(x) = g_2 \hat{n}_R(x) \hat{n}_L(x) + \frac{g_4}{2} (\hat{n}_R(x) \hat{n}_R(x) + \hat{n}_L(x) \hat{n}_L(x))$$



The diagram shows a horizontal line representing a 1D chain. A red arrow labeled 'R' points to the right, and a blue arrow labeled 'L' points to the left. The arrows are contained within a light orange shaded rectangular region.

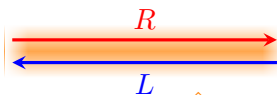
$$\hat{\mathcal{H}}_0(x) = iv_F \left(\hat{\psi}_R^\dagger(x) \partial_x \hat{\psi}_R(x) - \hat{\psi}_L^\dagger(x) \partial_x \hat{\psi}_L(x) \right)$$

$$\hat{\mathcal{H}}_{int}(x) = g_2 \hat{n}_R(x) \hat{n}_L(x) + \frac{g_4}{2} (\hat{n}_R(x) \hat{n}_R(x) + \hat{n}_L(x) \hat{n}_L(x))$$

Exact diagonalization in terms of **chiral boson fields**:

$$\hat{\mathcal{H}}_{LL}(x, t) = \frac{u}{2} \left\{ [\partial_x \hat{\phi}_+(x - ut)]^2 + [\partial_x \hat{\phi}_-(x + ut)]^2 \right\}$$

- Chiral boson fields are decoupled from each other.



$$\hat{\mathcal{H}}_0(x) = iv_F \left(\hat{\psi}_R^\dagger(x) \partial_x \hat{\psi}_R(x) - \hat{\psi}_L^\dagger(x) \partial_x \hat{\psi}_L(x) \right)$$

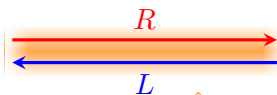
$$\hat{\mathcal{H}}_{int}(x) = g_2 \hat{n}_R(x) \hat{n}_L(x) + \frac{g_4}{2} (\hat{n}_R(x) \hat{n}_R(x) + \hat{n}_L(x) \hat{n}_L(x))$$

Exact diagonalization in terms of **chiral boson fields**:

$$\hat{\mathcal{H}}_{LL}(x, t) = \frac{u}{2} \left\{ [\partial_x \hat{\phi}_+(x - ut)]^2 + [\partial_x \hat{\phi}_-(x + ut)]^2 \right\}$$

$$K = \sqrt{\frac{2\pi v_F + g_4 - g_2}{2\pi v_F + g_4 + g_2}} \leq 1$$

- Chiral boson fields are decoupled from each other.
- Different interaction strength (K) correspond to different couple of chiral fields (Bogoliubov transformation).



$$\hat{\mathcal{H}}_0(x) = iv_F \left(\hat{\psi}_R^\dagger(x) \partial_x \hat{\psi}_R(x) - \hat{\psi}_L^\dagger(x) \partial_x \hat{\psi}_L(x) \right)$$

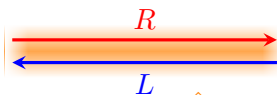
$$\hat{\mathcal{H}}_{int}(x) = g_2 \hat{n}_R(x) \hat{n}_L(x) + \frac{g_4}{2} (\hat{n}_R(x) \hat{n}_R(x) + \hat{n}_L(x) \hat{n}_L(x))$$

Exact diagonalization in terms of **chiral boson fields**:

$$\hat{\mathcal{H}}_{LL}(x, t) = \frac{u}{2} \left\{ [\partial_x \hat{\phi}_+(x - ut)]^2 + [\partial_x \hat{\phi}_-(x + ut)]^2 \right\} \quad \hat{\psi}_{R/L}^\dagger \propto e^{i\sqrt{2\pi}\hat{\phi}_{R/L}}$$

$$K = \sqrt{\frac{2\pi v_F + g_4 - g_2}{2\pi v_F + g_4 + g_2}} \leq 1 \quad \hat{\phi}_{R/L} = A_+(K) \hat{\phi}_\pm - A_-(K) \hat{\phi}_\mp$$

- Chiral boson fields are decoupled from each other.
- Different interaction strength (K) correspond to different couple of chiral fields (Bogoliubov transformation).



$$\hat{\mathcal{H}}_0(x) = iv_F \left(\hat{\psi}_R^\dagger(x) \partial_x \hat{\psi}_R(x) - \hat{\psi}_L^\dagger(x) \partial_x \hat{\psi}_L(x) \right)$$

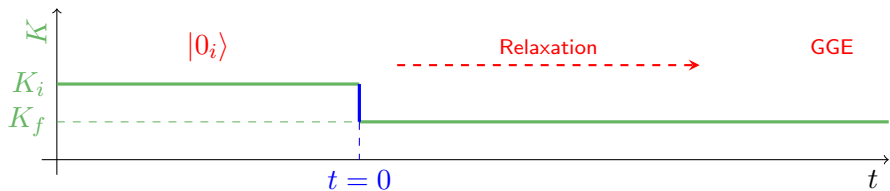
$$\hat{\mathcal{H}}_{int}(x) = g_2 \hat{n}_R(x) \hat{n}_L(x) + \frac{g_4}{2} (\hat{n}_R(x) \hat{n}_R(x) + \hat{n}_L(x) \hat{n}_L(x))$$

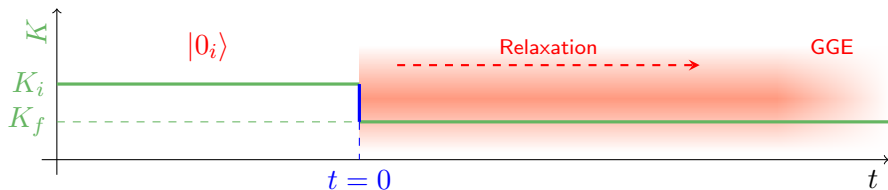
Exact diagonalization in terms of **chiral boson fields**:

$$\hat{\mathcal{H}}_{LL}(x, t) = \frac{u}{2} \left\{ [\partial_x \hat{\phi}_+(x - ut)]^2 + [\partial_x \hat{\phi}_-(x + ut)]^2 \right\} \quad \hat{\psi}_{R/L}^\dagger \propto e^{i\sqrt{2\pi}\hat{\phi}_{R/L}}$$

$$K = \sqrt{\frac{2\pi v_F + g_4 - g_2}{2\pi v_F + g_4 + g_2}} \leq 1 \quad \hat{\phi}_{R/L} = A_+(K) \hat{\phi}_\pm - A_-(K) \hat{\phi}_\mp$$

- Chiral boson fields are decoupled from each other.
- Different interaction strength (K) correspond to different couple of chiral fields (Bogoliubov transformation).
- LLs have peculiar features: zero bias anomaly in DOS (power-law behavior), fractionalization phenomena, ...





Focusing on the evolution **after** the quench:

$$\hat{\mathcal{H}}_{LL,f}(x, t > 0) = \frac{u}{2} \left\{ [\partial_x \hat{\phi}_{f,+}(x - ut)]^2 + [\partial_x \hat{\phi}_{f,-}(x + ut)]^2 \right\}$$

$$|0_i\rangle \propto \exp \left[\sigma[K_i, K_f] \int_{-\infty}^{+\infty} \left(\partial_x \hat{\phi}_{f,+}(x) \right) \hat{\phi}_{f,-}(x) \right] |0_f\rangle$$

Quench-induced entanglement between the two chiral channel $+$ and $-$.

Quench-induced
entanglement



Non-vanishing
local cross correlator
between bosonic fields

$$D_{cc}(t, \tau) = \langle \hat{\phi}_{+,f}(t - \tau) \hat{\phi}_{-,f}(t) \rangle + \langle \hat{\phi}_{+,f}(t) \hat{\phi}_{-,f}(t - \tau) \rangle \\ - \langle \hat{\phi}_{+,f}(t - \tau) \hat{\phi}_{-,f}(t - \tau) \rangle - \langle \hat{\phi}_{+,f}(t) \hat{\phi}_{-,f}(t) \rangle$$

Quench-induced entanglement \longrightarrow Non-vanishing local cross correlator between bosonic fields

$$D_{cc}(t, \tau) = \langle \hat{\phi}_{+,f}(t-\tau) \hat{\phi}_{-,f}(t) \rangle + \langle \hat{\phi}_{+,f}(t) \hat{\phi}_{-,f}(t-\tau) \rangle \\ - \langle \hat{\phi}_{+,f}(t-\tau) \hat{\phi}_{-,f}(t-\tau) \rangle - \langle \hat{\phi}_{+,f}(t) \hat{\phi}_{-,f}(t) \rangle$$

In the long time limit $t \gg \tau$, the cross correlator relaxes to zero as a power-law

$$D_{cc}(t, \tau) \approx \frac{d(\tau)}{t^2} \rightarrow 0$$

- Universal: the exponent is independent on both K_i and K_f
- Pure non-equilibrium effect
- Probe of quench-induced entanglement and relaxation dynamics

Local (lesser) non-equilibrium spectral function (NESF)

$$\mathcal{A}^<(\omega, t) = \frac{-i}{2\pi} \int e^{i\omega\tau} G^<(t, t-\tau) d\tau \quad G^<(t, t-\tau) = i\langle \hat{\psi}^\dagger(t-\tau) \hat{\psi}(\tau) \rangle$$

$$\mathcal{A}^<(\omega, t) - \mathcal{A}_\infty^<(\omega) \approx \frac{U_{\mathcal{A}}(\omega)}{t^2} + \frac{M_{\mathcal{A}}(\omega, t)}{t^\nu} \quad \text{with} \quad M_{\mathcal{A}} \approx e^{i\omega t}$$

$$\nu(K_i, K_f) \gtrsim 1$$

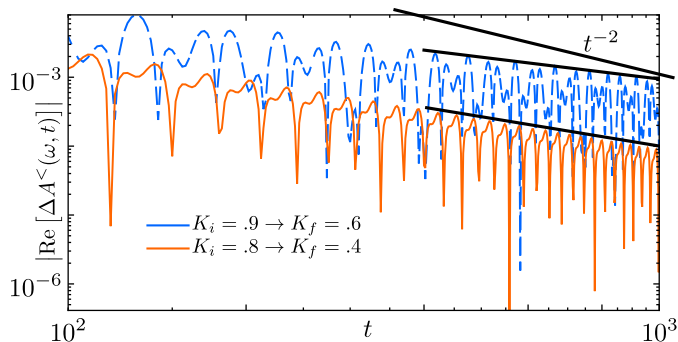
- universal power-law induced by the peculiar t dependence of D_{cc}
- non universal power-law (breaking of time translational invariance)

Local (lesser) non-equilibrium spectral function (NESF)

$$A^<(\omega, t) = \frac{-i}{2\pi} \int e^{i\omega\tau} G^<(t, t-\tau) d\tau \quad G^<(t, t-\tau) = i\langle \hat{\psi}^\dagger(t-\tau) \hat{\psi}(\tau) \rangle$$

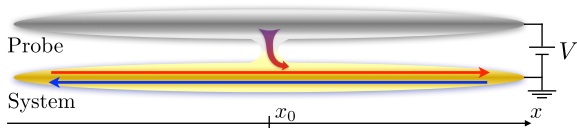
$$A^<(\omega, t) - A_\infty^<(\omega) \approx \frac{U_{\mathcal{A}}(\omega)}{t^2} + \frac{M_{\mathcal{A}}(\omega, t)}{t^\nu} \quad \text{with} \quad M_{\mathcal{A}} \approx e^{i\omega t}$$

$$\nu(K_i, K_f) \gtrsim 1$$



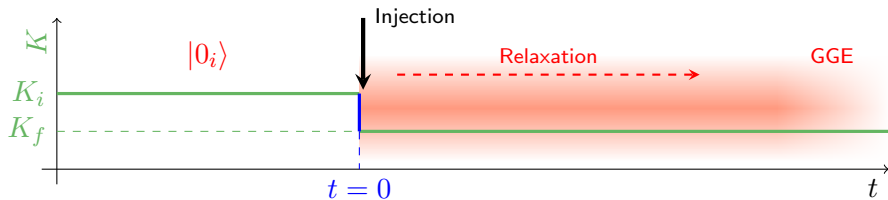
Although present,
the t^{-2} relaxation
is subleading!

Injection from a biased probe, locally tunnel coupled with the quenched system



$$\hat{\mathcal{H}}_p(x) = iv_F \hat{\chi}^\dagger(x) \partial_x \hat{\chi}(x)$$

$$\hat{H}_T(t) = \theta(t) \left[\lambda \hat{\psi}_R^\dagger(x_0) \hat{\chi}(x_0) + \text{h.c.} \right]$$



Breaking of inversion symmetry (injection of R electrons): allows to study also fractionalization of the injected currents.

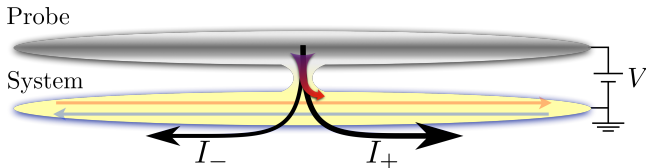
$$\hat{\psi}_R^\dagger \propto e^{i\sqrt{2\pi}(A_+ \hat{\phi}_+ + A_- \hat{\phi}_-)}$$

$$I_\eta(V, t) = \partial_t \int_{-\infty}^{\infty} \langle \delta n_\eta(V, x, t) \rangle dx \quad \text{Variation of chiral } (\eta = \pm) \text{ charge density (order } |\lambda|^2 \text{)}$$

$$= \left(\frac{1+\eta K_f}{2} \right) |\lambda|^2 2\text{Re} \left[\int_0^t i G^<(t, t-\tau) G_p^>(\tau) \sin(V\tau) d\tau \right]$$

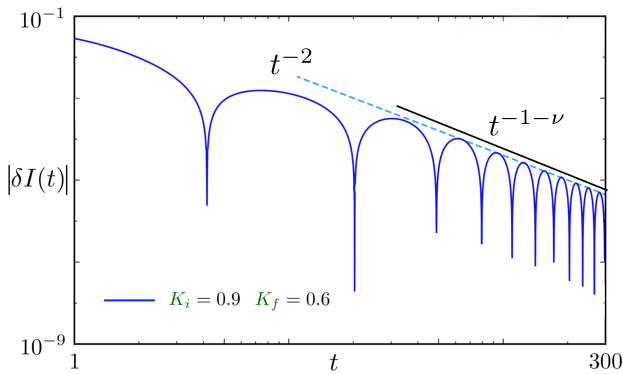
$$I_\eta(V, t) - I_\eta^\infty(V) \approx \left(\frac{1+\eta K_f}{2} \right) \left(\frac{U_I(V)}{t^2} + \frac{M_I(V, t)}{t^{1+\nu}} \right) \quad \nu(K_i, K_f) \gtrsim 1$$

$$M_I \approx \cos(Vt)$$



$$\begin{aligned}
 I_\eta(V, t) &= \partial_t \int_{-\infty}^{\infty} \langle \delta n_\eta(V, x, t) \rangle dx \\
 &= \left(\frac{1+\eta K_f}{2} \right) |\lambda|^2 2\text{Re} \left[\int_0^t i G^<(t, t-\tau) G_p^>(\tau) \sin(V\tau) d\tau \right]
 \end{aligned}$$

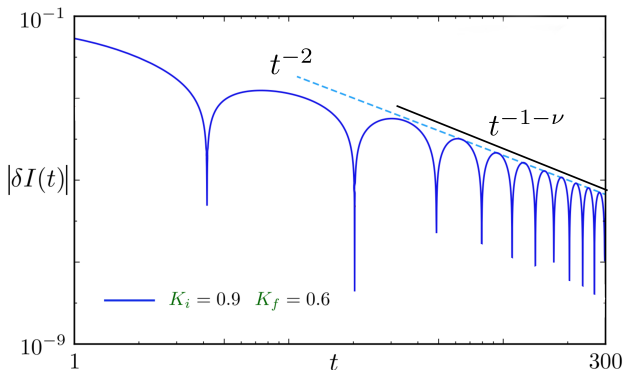
$$I_\eta(V, t) - I_\eta^\infty(V) \approx \left(\frac{1+\eta K_f}{2} \right) \left(\frac{U_I(V)}{t^2} + \frac{M_I(V, t)}{t^{1+\nu}} \right) \quad \begin{array}{l} \nu(K_i, K_f) \gtrsim 1 \\ M_I \approx \cos(Vt) \end{array}$$



- The **universal** decay is leading!
- Competition is still strong!

$$\begin{aligned}
 I_\eta(V, t) &= \partial_t \int_{-\infty}^{\infty} \langle \delta n_\eta(V, x, t) \rangle dx \\
 &= \left(\frac{1+\eta K_f}{2} \right) |\lambda|^2 2\text{Re} \left[\int_0^t i G^<(t, t-\tau) G_p^>(\tau) \sin(V\tau) d\tau \right]
 \end{aligned}$$

$$I_\eta(V, t) - I_\eta^\infty(V) \approx \left(\frac{1+\eta K_f}{2} \right) \left(\frac{U_I(V)}{t^2} + \frac{M_I(V, t)}{t^{1+\nu}} \right) \quad \begin{array}{l} \nu(K_i, K_f) \gtrsim 1 \\ M_I \approx \cos(Vt) \end{array}$$



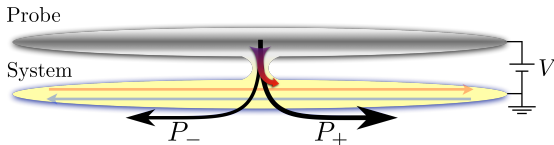
- The **universal** decay is leading!
- Competition is still strong!
- **Charge fractionalization** is “trivial”!

$$P_\eta(V, t) = \partial_t \int_{-\infty}^{\infty} \langle \delta \mathcal{H}_\eta(V, x, t) \rangle dx \quad \text{Variation of chiral } (\eta = \pm) \text{ hamiltonian density (order } |\lambda|^2)$$

$$= 2|\lambda|^2 \text{Im} \left[\int_0^t G_P^>(\tau) (\partial_{\bar{t}} - \eta u_f \partial_\xi) G_R^<(\xi; \bar{t}, t - \tau) \Big|_{\substack{\xi=0 \\ \bar{t}=t}} \cos(V\tau) d\tau \right]$$

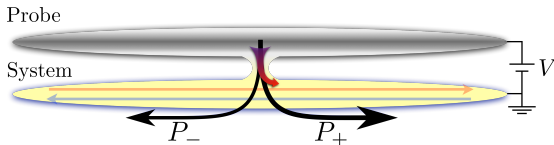
$$P_\eta(V, t) - P_\eta^\infty(V) \approx \mathcal{R}_\eta(K_f) \left(\frac{U_P^A(V)}{t^2} + \frac{M_P^A(V, t)}{t^{2+\nu}} \right) \quad \nu(K_i, K_f) \gtrsim 1$$

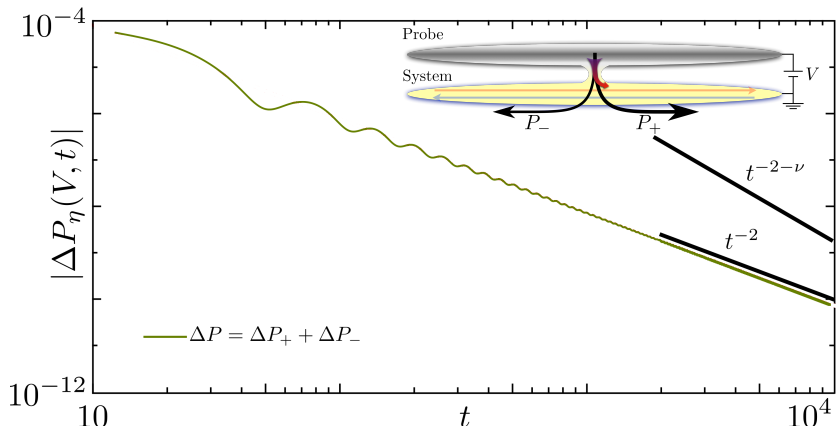
$$+ \xi(K_i, K_f) \left(\frac{U_P^B(V)}{t^2} + \frac{M_P^B(V, t)}{t^{2+\nu}} \right) \quad M_P \approx \cos(Vt)$$



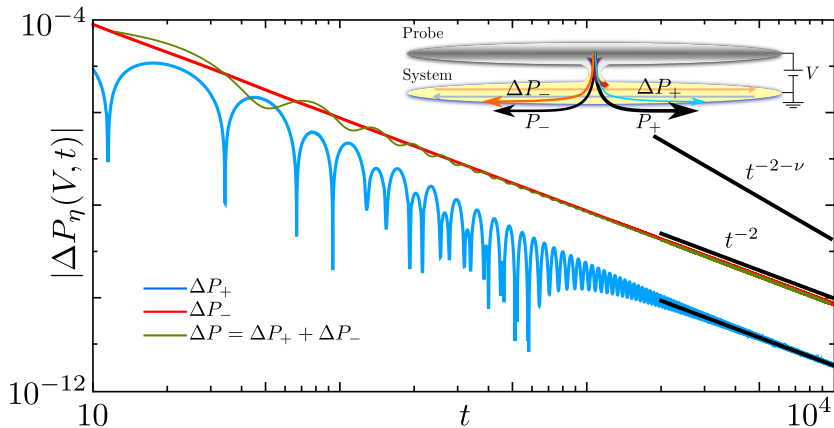
$$\begin{aligned}
 P_\eta(V, t) &= \partial_t \int_{-\infty}^{\infty} \langle \delta \mathcal{H}_\eta(V, x, t) \rangle dx && \text{Variation of chiral } (\eta = \pm) \\
 & && \text{hamiltonian density (order } |\lambda|^2) \\
 &= 2|\lambda|^2 \text{Im} \left[\int_0^t G_P^>(\tau) (\partial_{\bar{t}} - \eta u_f \partial_\xi) G_R^<(\xi; \bar{t}, t - \tau) \Big|_{\substack{\xi=0 \\ \bar{t}=t}} \cos(V\tau) d\tau \right]
 \end{aligned}$$

$$\begin{aligned}
 P_\eta(V, t) - P_\eta^\infty(V) &\approx \mathcal{R}_\eta(K_f) \left(\frac{U_P^A(V)}{t^2} + \frac{M_P^A(V, t)}{t^{2+\nu}} \right) && \nu(K_i, K_f) \gtrsim 1 \\
 &+ \xi(K_i, K_f) \left(\frac{U_P^B(V)}{t^2} + \frac{M_P^B(V, t)}{t^{2+\nu}} \right) && M_P \approx \cos(Vt)
 \end{aligned}$$





$$K_i = 0.9 \quad \rightarrow \quad K_f = 0.6$$



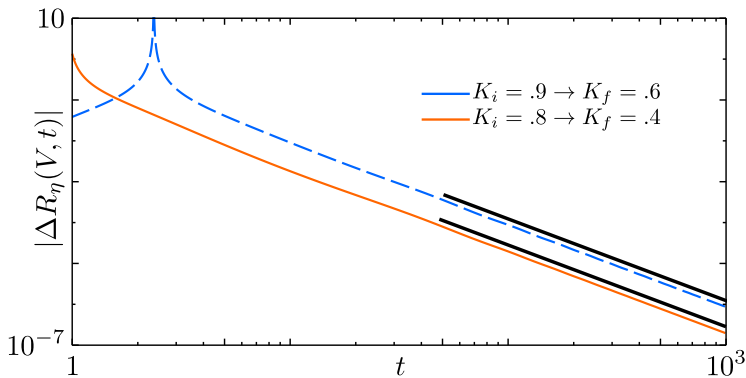
$$K_i = 0.9 \quad \rightarrow \quad K_f = 0.6$$

The **left moving component** of the transient (ΔP_-) displays an almost perfect t^{-2} behavior, even after a short time.

Fractionalization ratio

$$R_+(V, t) = \frac{P_+(V, t)}{P_+(V, t) + P_-(V, t)}$$

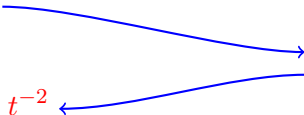
$$\Delta R_+(V, t) = R_+(V, t) - R_+^\infty(t)$$



Interaction quench
in Luttinger liquid

Non-equilibrium
entangled state

Universal relaxation $\propto t^{-2}$
of local cross-correlators

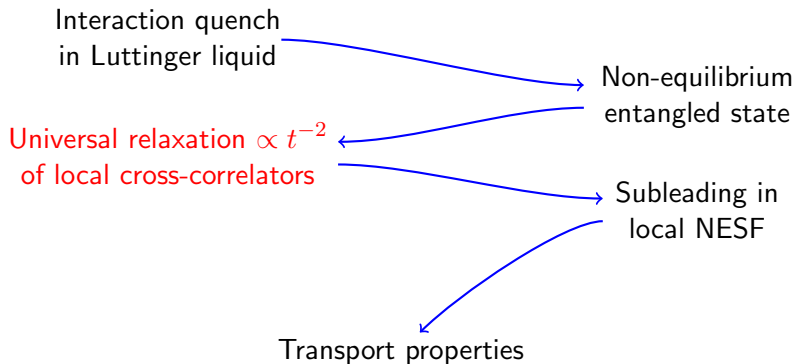


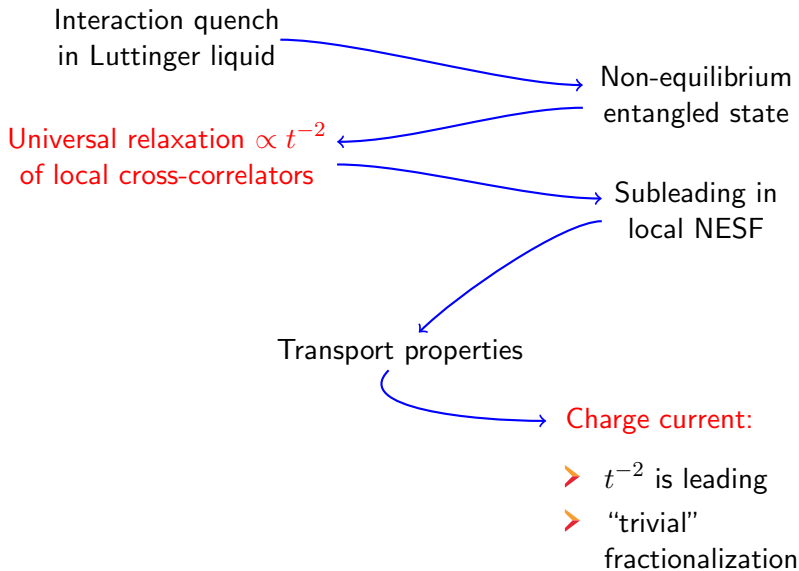
Interaction quench
in Luttinger liquid

Non-equilibrium
entangled state

Universal relaxation $\propto t^{-2}$
of local cross-correlators

Subleading in
local NESF





Interaction quench
in Luttinger liquid

Non-equilibrium
entangled state

Universal relaxation $\propto t^{-2}$
of local cross-correlators

Subleading in
local NESF

Transport properties

Energy current:

Charge current:

- t^{-2} emerges even more clearly
- non-trivial partitioning:
 - Predominance of the “-” channel
 - Direct fingerprint of quench-induced dynamics

- t^{-2} is leading
- “trivial” fractionalization

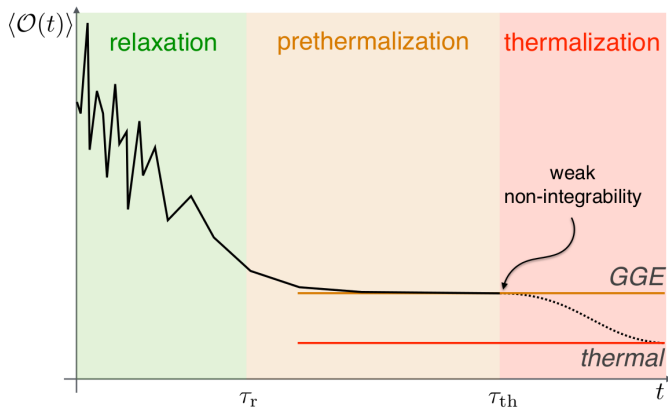
A. Calzona, F. M. Gambetta, F. Cavaliere, M. Carrega, and M. Sassetti
arXiv:1706.01676

Thank you for your attention!

$$\nu = \frac{K_f^4 + K_i^2 + 3K_f^2(1 + K_i^2)}{8K_f^2K_i}$$

Integrable system \rightarrow GGE

Non-integrable system (ETH) \rightarrow Thermalization



Polkovnikov, RMP (2012)
 D'Alessio, Adv. Phys (2016)
 Rigol, PRL (2007)

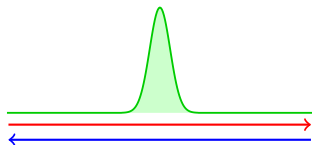
- Breakdown of Fermi liquid model, **Luttinger liquid** instead!
- Excitations are **collective** and with **bosonic** nature

- Breakdown of Fermi liquid model, **Luttinger liquid** instead!
- Excitations are **collective** and with **bosonic** nature

$$\hat{\mathcal{H}}_0(x) = iv_F \left(\hat{\psi}_R^\dagger(x) \partial_x \hat{\psi}_R(x) - \hat{\psi}_L^\dagger(x) \partial_x \hat{\psi}_L(x) \right)$$

Injecting a **R** electron:

$$\hat{\psi}_R^\dagger(x, t) = \hat{\psi}_R^\dagger(x - v_F t)$$

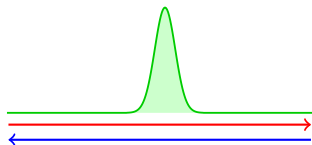


- Breakdown of Fermi liquid model, **Luttinger liquid** instead!
- Excitations are **collective** and with **bosonic** nature

$$\hat{\mathcal{H}}_0(x) = iv_F \left(\hat{\psi}_R^\dagger(x) \partial_x \hat{\psi}_R(x) - \hat{\psi}_L^\dagger(x) \partial_x \hat{\psi}_L(x) \right)$$

Injecting a **R** electron:

$$\hat{\psi}_R^\dagger(x, t) = \hat{\psi}_R^\dagger(x - v_F t)$$

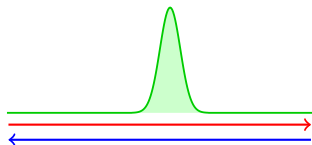


- Breakdown of Fermi liquid model, **Luttinger liquid** instead!
- Excitations are **collective** and with **bosonic** nature

$$\hat{\mathcal{H}}_0(x) = iv_F \left(\hat{\psi}_R^\dagger(x) \partial_x \hat{\psi}_R(x) - \hat{\psi}_L^\dagger(x) \partial_x \hat{\psi}_L(x) \right)$$

Injecting a **R** electron:

$$\hat{\psi}_R^\dagger(x, t) = \hat{\psi}_R^\dagger(x - v_F t)$$

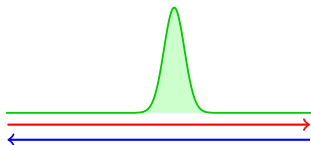


- Breakdown of Fermi liquid model, **Luttinger liquid** instead!
- Excitations are **collective** and with **bosonic** nature

$$\hat{\mathcal{H}}_0(x) = iv_F \left(\hat{\psi}_R^\dagger(x) \partial_x \hat{\psi}_R(x) - \hat{\psi}_L^\dagger(x) \partial_x \hat{\psi}_L(x) \right)$$

Injecting a **R** electron:

$$\hat{\psi}_R^\dagger(x, t) = \hat{\psi}_R^\dagger(x - v_F t)$$

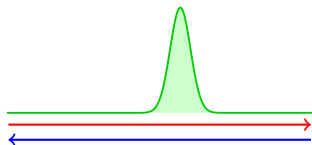


- Breakdown of Fermi liquid model, **Luttinger liquid** instead!
- Excitations are **collective** and with **bosonic** nature

$$\hat{\mathcal{H}}_0(x) = iv_F \left(\hat{\psi}_R^\dagger(x) \partial_x \hat{\psi}_R(x) - \hat{\psi}_L^\dagger(x) \partial_x \hat{\psi}_L(x) \right)$$

Injecting a **R** electron:

$$\hat{\psi}_R^\dagger(x, t) = \hat{\psi}_R^\dagger(x - v_F t)$$

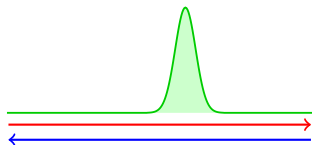


- Breakdown of Fermi liquid model, **Luttinger liquid** instead!
- Excitations are **collective** and with **bosonic** nature

$$\hat{\mathcal{H}}_0(x) = iv_F \left(\hat{\psi}_R^\dagger(x) \partial_x \hat{\psi}_R(x) - \hat{\psi}_L^\dagger(x) \partial_x \hat{\psi}_L(x) \right)$$

Injecting a **R** electron:

$$\hat{\psi}_R^\dagger(x, t) = \hat{\psi}_R^\dagger(x - v_F t)$$

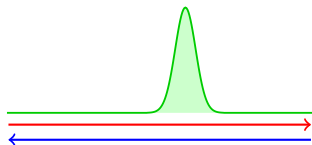


- Breakdown of Fermi liquid model, **Luttinger liquid** instead!
- Excitations are **collective** and with **bosonic** nature

$$\hat{\mathcal{H}}_0(x) = iv_F \left(\hat{\psi}_R^\dagger(x) \partial_x \hat{\psi}_R(x) - \hat{\psi}_L^\dagger(x) \partial_x \hat{\psi}_L(x) \right)$$

Injecting a **R** electron:

$$\hat{\psi}_R^\dagger(x, t) = \hat{\psi}_R^\dagger(x - v_F t)$$



- Breakdown of Fermi liquid model, **Luttinger liquid** instead!
- Excitations are **collective** and with **bosonic** nature

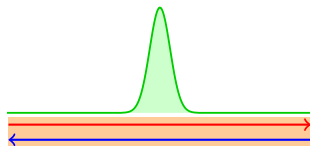
$$\hat{\mathcal{H}}_0(x) = iv_F \left(\hat{\psi}_R^\dagger(x) \partial_x \hat{\psi}_R(x) - \hat{\psi}_L^\dagger(x) \partial_x \hat{\psi}_L(x) \right)$$

$$K = \sqrt{\frac{2\pi v_F + g_4 - g_2}{2\pi v_F + g_4 + g_2}} \leq 1 \quad \hat{\mathcal{H}}_{int}(x) = g_2 \hat{n}_R(x) \hat{n}_L(x) + g_4 (\hat{n}_R(x) \hat{n}_R(x) + \hat{n}_L(x) \hat{n}_L(x)) / 2$$

$$\hat{\mathcal{H}}_{LL}(x, t) = \frac{u}{2} \left(\partial_x \hat{\phi}_+(x - ut) \right)^2 + \frac{u}{2} \left(\partial_x \hat{\phi}_-(x + ut) \right)^2$$

Injecting a **R** electron:

$$\hat{\psi}_R^\dagger(x, t) \simeq e^{i\sqrt{2\pi}(A_+ \hat{\phi}_+(x-ut) + A_- \hat{\phi}_-(x+ut))}$$



- Breakdown of Fermi liquid model, **Luttinger liquid** instead!
- Excitations are **collective** and with **bosonic** nature

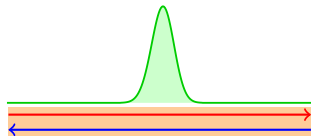
$$\hat{\mathcal{H}}_0(x) = iv_F \left(\hat{\psi}_R^\dagger(x) \partial_x \hat{\psi}_R(x) - \hat{\psi}_L^\dagger(x) \partial_x \hat{\psi}_L(x) \right)$$

$$K = \sqrt{\frac{2\pi v_F + g_4 - g_2}{2\pi v_F + g_4 + g_2}} \leq 1 \quad \hat{\mathcal{H}}_{int}(x) = g_2 \hat{n}_R(x) \hat{n}_L(x) + g_4 (\hat{n}_R(x) \hat{n}_R(x) + \hat{n}_L(x) \hat{n}_L(x)) / 2$$

$$\hat{\mathcal{H}}_{LL}(x, t) = \frac{u}{2} \left(\partial_x \hat{\phi}_+(x - ut) \right)^2 + \frac{u}{2} \left(\partial_x \hat{\phi}_-(x + ut) \right)^2$$

Injecting a **R** electron:

$$\hat{\psi}_R^\dagger(x, t) \simeq e^{i\sqrt{2\pi}(A_+ \hat{\phi}_+(x-ut) + A_- \hat{\phi}_-(x+ut))}$$



- Breakdown of Fermi liquid model, **Luttinger liquid** instead!
- Excitations are **collective** and with **bosonic** nature

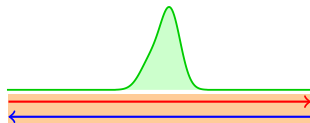
$$\hat{\mathcal{H}}_0(x) = iv_F \left(\hat{\psi}_R^\dagger(x) \partial_x \hat{\psi}_R(x) - \hat{\psi}_L^\dagger(x) \partial_x \hat{\psi}_L(x) \right)$$

$$K = \sqrt{\frac{2\pi v_F + g_4 - g_2}{2\pi v_F + g_4 + g_2}} \leq 1 \quad \hat{\mathcal{H}}_{int}(x) = g_2 \hat{n}_R(x) \hat{n}_L(x) + g_4 (\hat{n}_R(x) \hat{n}_R(x) + \hat{n}_L(x) \hat{n}_L(x)) / 2$$

$$\hat{\mathcal{H}}_{LL}(x, t) = \frac{u}{2} \left(\partial_x \hat{\phi}_+(x - ut) \right)^2 + \frac{u}{2} \left(\partial_x \hat{\phi}_-(x + ut) \right)^2$$

Injecting a **R** electron:

$$\hat{\psi}_R^\dagger(x, t) \simeq e^{i\sqrt{2\pi}(A_+ \hat{\phi}_+(x-ut) + A_- \hat{\phi}_-(x+ut))}$$



- Breakdown of Fermi liquid model, **Luttinger liquid** instead!
- Excitations are **collective** and with **bosonic** nature

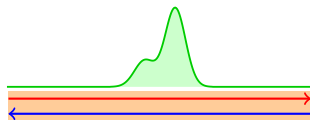
$$\hat{\mathcal{H}}_0(x) = iv_F \left(\hat{\psi}_R^\dagger(x) \partial_x \hat{\psi}_R(x) - \hat{\psi}_L^\dagger(x) \partial_x \hat{\psi}_L(x) \right)$$

$$K = \sqrt{\frac{2\pi v_F + g_4 - g_2}{2\pi v_F + g_4 + g_2}} \leq 1 \quad \hat{\mathcal{H}}_{int}(x) = g_2 \hat{n}_R(x) \hat{n}_L(x) + g_4 (\hat{n}_R(x) \hat{n}_R(x) + \hat{n}_L(x) \hat{n}_L(x)) / 2$$

$$\hat{\mathcal{H}}_{LL}(x, t) = \frac{u}{2} \left(\partial_x \hat{\phi}_+(x - ut) \right)^2 + \frac{u}{2} \left(\partial_x \hat{\phi}_-(x + ut) \right)^2$$

Injecting a **R** electron:

$$\hat{\psi}_R^\dagger(x, t) \simeq e^{i\sqrt{2\pi}(A_+ \hat{\phi}_+(x-ut) + A_- \hat{\phi}_-(x+ut))}$$



- Breakdown of Fermi liquid model, **Luttinger liquid** instead!
- Excitations are **collective** and with **bosonic** nature

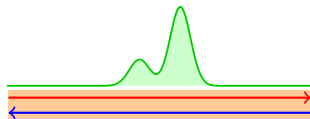
$$\hat{\mathcal{H}}_0(x) = iv_F \left(\hat{\psi}_R^\dagger(x) \partial_x \hat{\psi}_R(x) - \hat{\psi}_L^\dagger(x) \partial_x \hat{\psi}_L(x) \right)$$

$$K = \sqrt{\frac{2\pi v_F + g_4 - g_2}{2\pi v_F + g_4 + g_2}} \leq 1 \quad \hat{\mathcal{H}}_{int}(x) = g_2 \hat{n}_R(x) \hat{n}_L(x) + g_4 (\hat{n}_R(x) \hat{n}_R(x) + \hat{n}_L(x) \hat{n}_L(x)) / 2$$

$$\hat{\mathcal{H}}_{LL}(x, t) = \frac{u}{2} \left(\partial_x \hat{\phi}_+(x - ut) \right)^2 + \frac{u}{2} \left(\partial_x \hat{\phi}_-(x + ut) \right)^2$$

Injecting a **R** electron:

$$\hat{\psi}_R^\dagger(x, t) \simeq e^{i\sqrt{2\pi}(A_+ \hat{\phi}_+(x-ut) + A_- \hat{\phi}_-(x+ut))}$$



- Breakdown of Fermi liquid model, **Luttinger liquid** instead!
- Excitations are **collective** and with **bosonic** nature

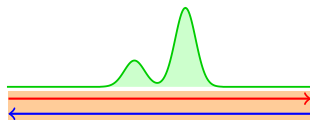
$$\hat{\mathcal{H}}_0(x) = iv_F \left(\hat{\psi}_R^\dagger(x) \partial_x \hat{\psi}_R(x) - \hat{\psi}_L^\dagger(x) \partial_x \hat{\psi}_L(x) \right)$$

$$K = \sqrt{\frac{2\pi v_F + g_4 - g_2}{2\pi v_F + g_4 + g_2}} \leq 1 \quad \hat{\mathcal{H}}_{int}(x) = g_2 \hat{n}_R(x) \hat{n}_L(x) + g_4 (\hat{n}_R(x) \hat{n}_R(x) + \hat{n}_L(x) \hat{n}_L(x)) / 2$$

$$\hat{\mathcal{H}}_{LL}(x, t) = \frac{u}{2} \left(\partial_x \hat{\phi}_+(x - ut) \right)^2 + \frac{u}{2} \left(\partial_x \hat{\phi}_-(x + ut) \right)^2$$

Injecting a **R** electron:

$$\hat{\psi}_R^\dagger(x, t) \simeq e^{i\sqrt{2\pi}(A_+ \hat{\phi}_+(x-ut) + A_- \hat{\phi}_-(x+ut))}$$



- Breakdown of Fermi liquid model, **Luttinger liquid** instead!
- Excitations are **collective** and with **bosonic** nature

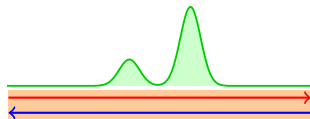
$$\hat{\mathcal{H}}_0(x) = iv_F \left(\hat{\psi}_R^\dagger(x) \partial_x \hat{\psi}_R(x) - \hat{\psi}_L^\dagger(x) \partial_x \hat{\psi}_L(x) \right)$$

$$K = \sqrt{\frac{2\pi v_F + g_4 - g_2}{2\pi v_F + g_4 + g_2}} \leq 1 \quad \hat{\mathcal{H}}_{int}(x) = g_2 \hat{n}_R(x) \hat{n}_L(x) + g_4 (\hat{n}_R(x) \hat{n}_R(x) + \hat{n}_L(x) \hat{n}_L(x)) / 2$$

$$\hat{\mathcal{H}}_{LL}(x, t) = \frac{u}{2} \left(\partial_x \hat{\phi}_+(x - ut) \right)^2 + \frac{u}{2} \left(\partial_x \hat{\phi}_-(x + ut) \right)^2$$

Injecting a **R** electron:

$$\hat{\psi}_R^\dagger(x, t) \simeq e^{i\sqrt{2\pi}(A_+ \hat{\phi}_+(x-ut) + A_- \hat{\phi}_-(x+ut))}$$



- Breakdown of Fermi liquid model, **Luttinger liquid** instead!
- Excitations are **collective** and with **bosonic** nature

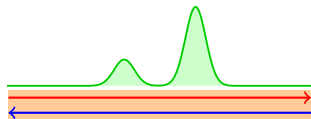
$$\hat{\mathcal{H}}_0(x) = iv_F \left(\hat{\psi}_R^\dagger(x) \partial_x \hat{\psi}_R(x) - \hat{\psi}_L^\dagger(x) \partial_x \hat{\psi}_L(x) \right)$$

$$K = \sqrt{\frac{2\pi v_F + g_4 - g_2}{2\pi v_F + g_4 + g_2}} \leq 1 \quad \hat{\mathcal{H}}_{int}(x) = g_2 \hat{n}_R(x) \hat{n}_L(x) + g_4 (\hat{n}_R(x) \hat{n}_R(x) + \hat{n}_L(x) \hat{n}_L(x)) / 2$$

$$\hat{\mathcal{H}}_{LL}(x, t) = \frac{u}{2} \left(\partial_x \hat{\phi}_+(x - ut) \right)^2 + \frac{u}{2} \left(\partial_x \hat{\phi}_-(x + ut) \right)^2$$

Injecting a **R** electron:

$$\hat{\psi}_R^\dagger(x, t) \simeq e^{i\sqrt{2\pi}(A_+ \hat{\phi}_+(x-ut) + A_- \hat{\phi}_-(x+ut))}$$



- Breakdown of Fermi liquid model, **Luttinger liquid** instead!
- Excitations are **collective** and with **bosonic** nature

$$\hat{\mathcal{H}}_0(x) = iv_F \left(\hat{\psi}_R^\dagger(x) \partial_x \hat{\psi}_R(x) - \hat{\psi}_L^\dagger(x) \partial_x \hat{\psi}_L(x) \right)$$

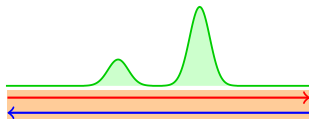
$$K = \sqrt{\frac{2\pi v_F + g_4 - g_2}{2\pi v_F + g_4 + g_2}} \leq 1 \quad \hat{\mathcal{H}}_{int}(x) = g_2 \hat{n}_R(x) \hat{n}_L(x) + g_4 (\hat{n}_R(x) \hat{n}_R(x) + \hat{n}_L(x) \hat{n}_L(x)) / 2$$

$$\hat{\mathcal{H}}_{LL}(x, t) = \frac{u}{2} \left(\partial_x \hat{\phi}_+(x - ut) \right)^2 + \frac{u}{2} \left(\partial_x \hat{\phi}_-(x + ut) \right)^2$$

Injecting a **R** electron:

$$\hat{\psi}_R^\dagger(x, t) \simeq e^{i\sqrt{2\pi}(A_+ \hat{\phi}_+(x-ut) + A_- \hat{\phi}_-(x+ut))}$$

Fractional excitations!!



- Breakdown of Fermi liquid model, **Luttinger liquid** instead!
- Excitations are **collective** and with **bosonic** nature

$$\hat{\mathcal{H}}_0(x) = iv_F \left(\hat{\psi}_R^\dagger(x) \partial_x \hat{\psi}_R(x) - \hat{\psi}_L^\dagger(x) \partial_x \hat{\psi}_L(x) \right)$$

$$K = \sqrt{\frac{2\pi v_F + g_4 - g_2}{2\pi v_F + g_4 + g_2}} \leq 1 \quad \hat{\mathcal{H}}_{int}(x) = g_2 \hat{n}_R(x) \hat{n}_L(x) + g_4 (\hat{n}_R(x) \hat{n}_R(x) + \hat{n}_L(x) \hat{n}_L(x)) / 2$$

$$\hat{\mathcal{H}}_{LL}(x, t) = \frac{u}{2} \left(\partial_x \hat{\phi}_+(x - ut) \right)^2 + \frac{u}{2} \left(\partial_x \hat{\phi}_-(x + ut) \right)^2$$

Injecting a **R** electron:

$$\hat{\psi}_R^\dagger(x, t) \simeq e^{i\sqrt{2\pi}(A_+ \hat{\phi}_+(x-ut) + A_- \hat{\phi}_-(x+ut))}$$

Fractional excitations!!

