

Optimal quantum interference thermoelectric heat engine with edge states

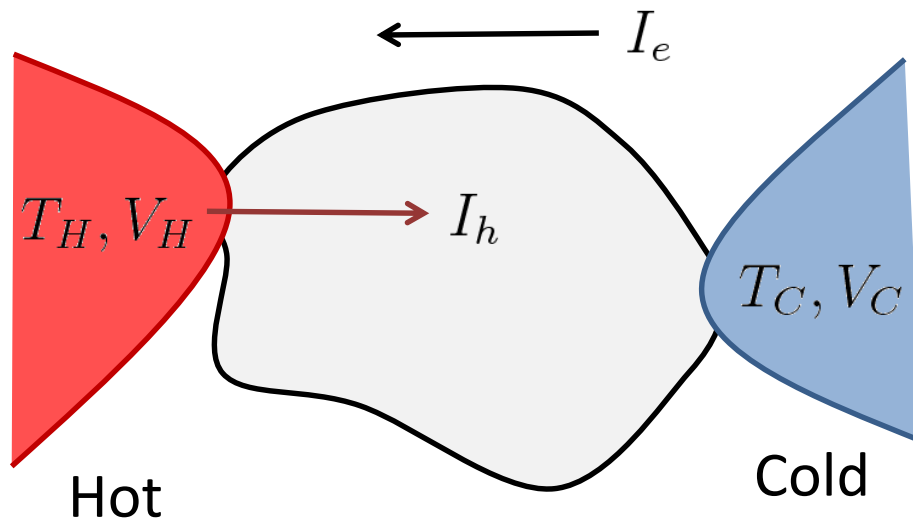
Peter Samuelsson, Sara Kheradsoud, Björn Sothmann

Mesoscopic Transport and Quantum Coherence
Aalto University, August 2017

P. Samuelsson, S. Kheradsoud, B. Sothmann, Phys. Rev. Lett. 118, 256801 (2017) .

Mesoscopic thermoelectric heat engine

Steady state heat to electrical work conversion in mesoscopic conductors



Motivation

- Proof-of-principle for nanoscale waste heat recovery.
- Role of coherence in heat engine performance.
- Investigate fundamentals of energy transport.

Performance

- Power

$$P = I_e V$$

Electrical current I_e flowing against applied bias

$$V = V_H - V_C > 0$$

- Efficiency

$$\eta = \frac{P}{I_h} \leq \eta_C$$

bounded by Carnot η_C .

Optimal performance

Two-terminal heat engine, linear response

Transmission

- Energy dependence of

$$T(\varepsilon)$$

governs thermoelectric properties

- Maximizing output power

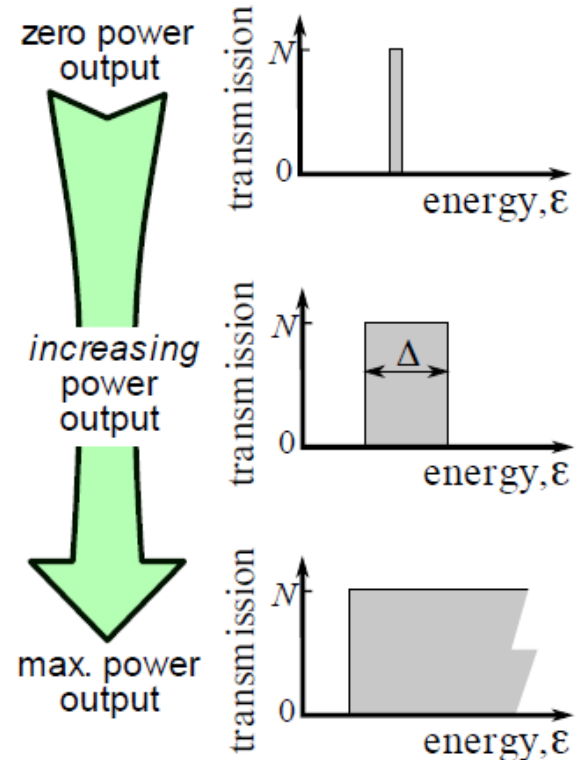


Step function in energy

Q1: Can a purely interference based TE-heat engine be optimal?

Optimization

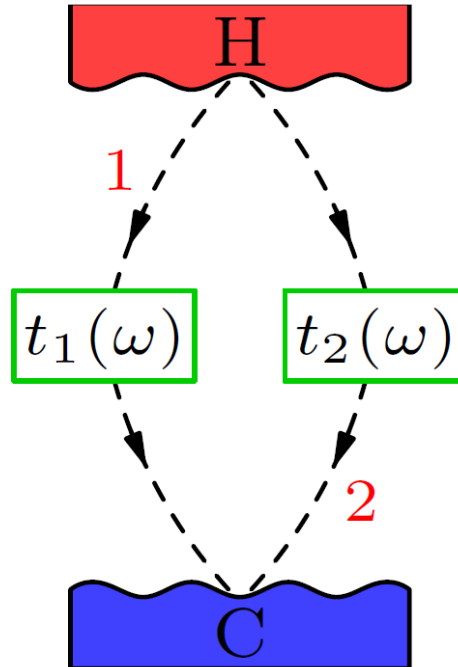
Transmission giving optimal η for given P .



Whitney, PRL 2014

Double slit interferometer

Generic interferometer



Scattering amplitude

$$A(\omega) \propto t_1(\omega) + e^{i\phi_0} t_2(\omega)$$

ϕ_0

Phase due to
Aharonov-Bohm, ..

No individual TE-effect

$$T_i = |t_i(\omega)|^2$$

Energy
independent

$$\Rightarrow t_i(\omega) = \sqrt{T_i} e^{i\alpha_i(\omega)}$$

Energy
dependence

Scattering probability

$$|A(\omega)|^2 \propto T_1 + T_2 + 2\sqrt{T_1 T_2} \cos[\alpha_1(\omega) - \alpha_2(\omega) - \phi_0]$$

Double slit interferometer

$$|A(\omega)|^2 \propto T_1 + T_2 + 2\sqrt{T_1 T_2} \cos[\alpha_1(\omega) - \alpha_2(\omega) - \phi_0]$$

Conditions for transmission step

- Symmetric interferometer, $T_1 = T_2$.
- Sharp jump $\pi \rightarrow 0$ for $\alpha_1(\omega) - \alpha_2(\omega) - \phi_0$ at $\omega = \omega_0$.



$$T(\omega) = \theta(\omega - \omega_0)$$

Step function
transmission

Q1: Can a purely interference based TE-heat engine be optimal?

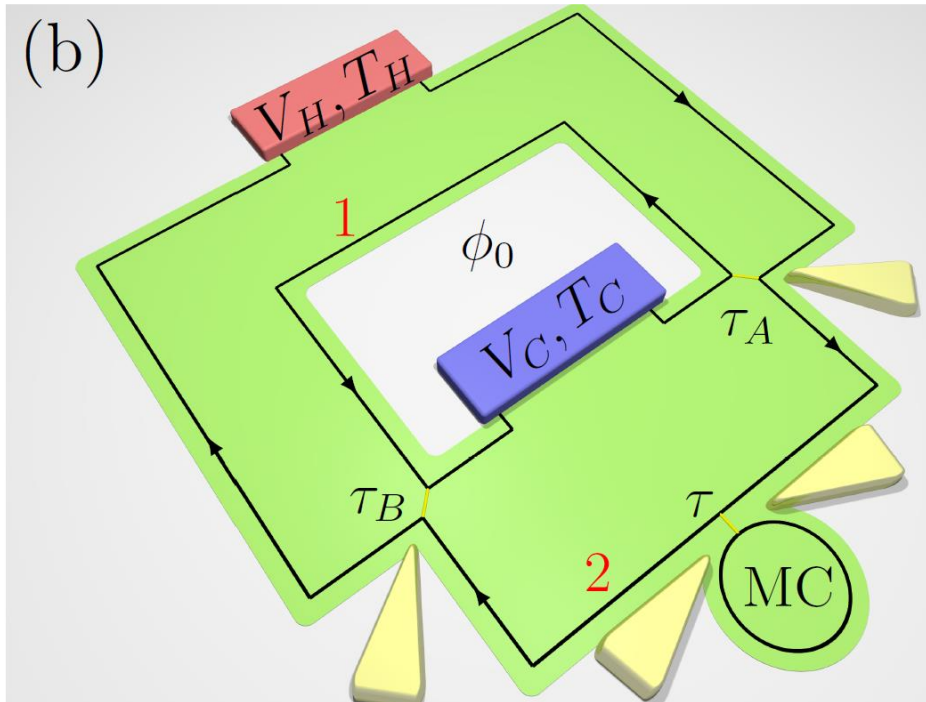
A1: Yes



Q2: Is there a possible experimental realization of such optimal engine?

Edge state Mach-Zehnder with capacitor

Extending Hofer, Sothmann, PRB 2015



Properties

- Quantum point contacts
 τ_A, τ_B
- Mesoscopic capacitor (MC)
 τ, Δ, ω_0 Fève et al, Science 2007
- Contacts
 $T_C, T_H \quad V_C, V_H$
- Interferometer
 ϕ_0

Total transmission amplitude (equal arm lengths)

$$t(\omega) = \sqrt{\tau_A \tau_B} - \sqrt{(1 - \tau_A)(1 - \tau_B)} e^{i[\alpha(\omega) + \phi_0]}$$

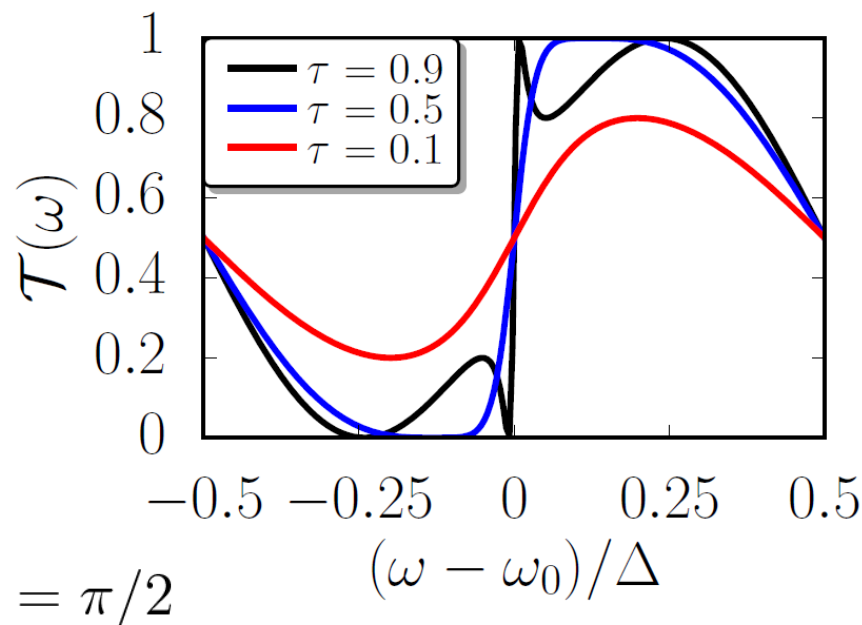
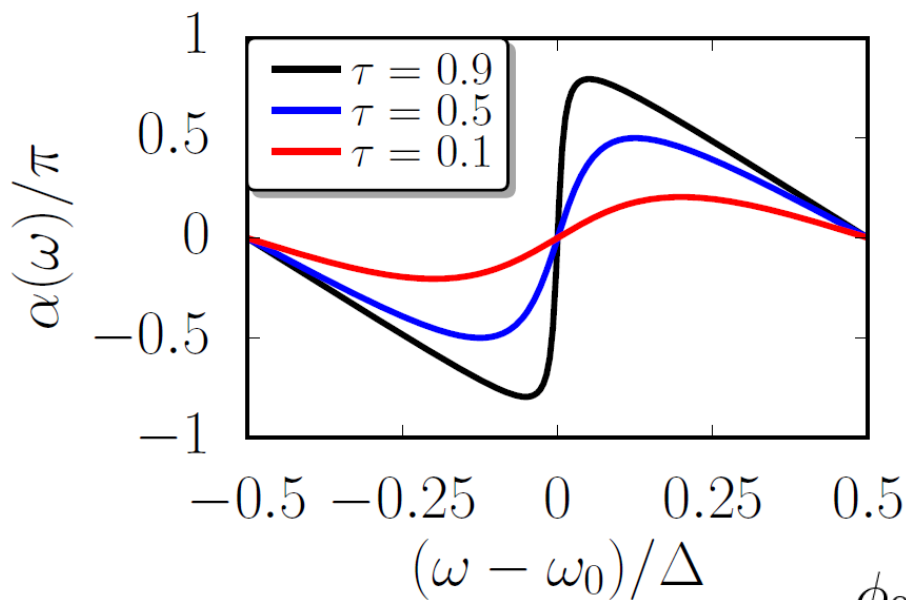
$$\alpha(\omega) = 2 \arctan \frac{\sqrt{\tau} \sin \left(2\pi \frac{\omega - \omega_0}{\Delta} \right)}{1 - \sqrt{\tau} \cos \left(2\pi \frac{\omega - \omega_0}{\Delta} \right)}$$

Scattering
phase at MC

Transmission properties

Semitransparent splitters $\tau_A = \tau_B = 1/2$ (symmetry condition),

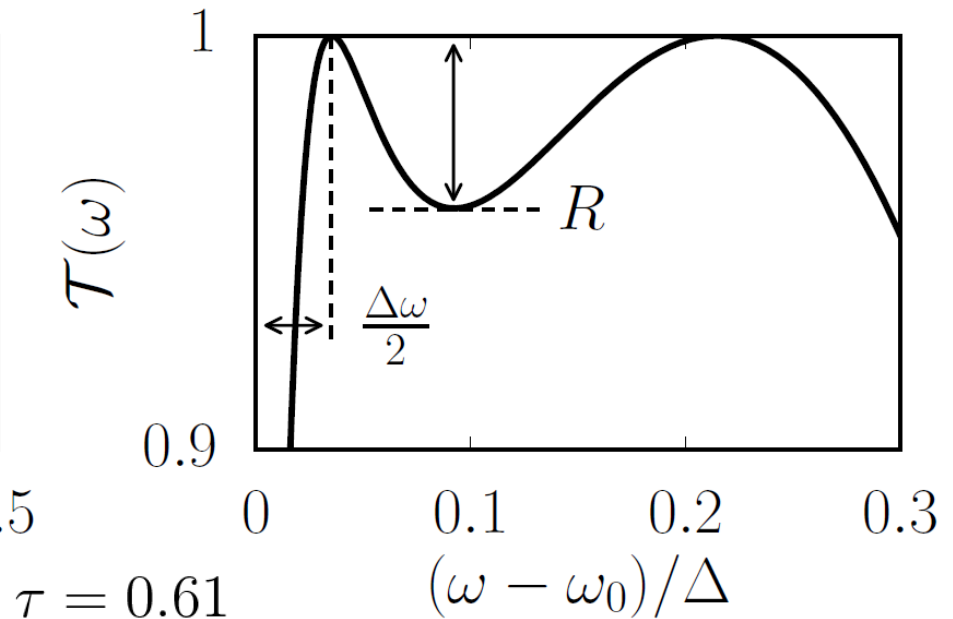
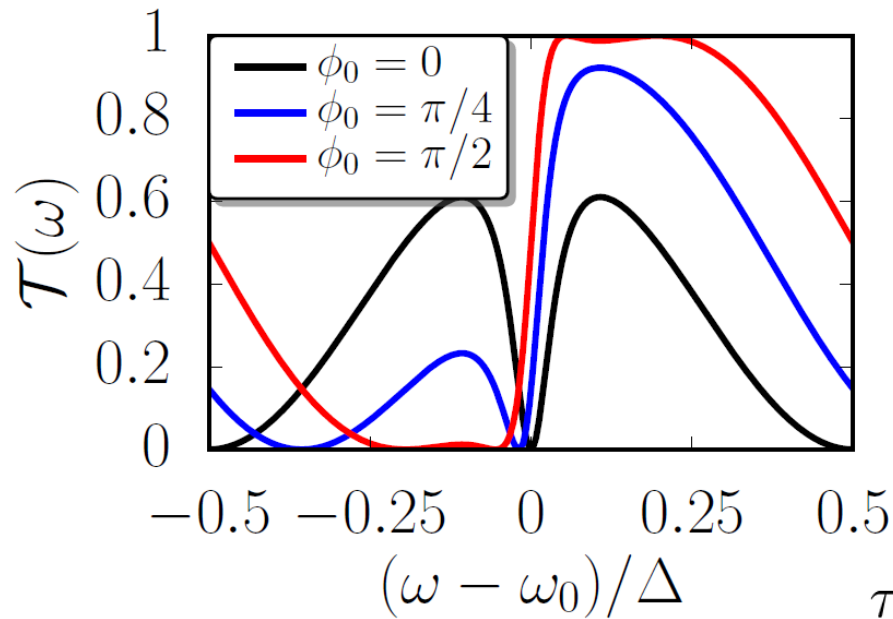
$$\mathcal{T}(\omega) = \frac{\left[\sin\left(\frac{\phi_0}{2}\right) - \sqrt{\tau} \sin\left(\frac{\phi_0}{2} - 2\pi\frac{\omega - \omega_0}{\Delta}\right) \right]^2}{1 - 2\sqrt{\tau} \cos\left(2\pi\frac{\omega - \omega_0}{\Delta}\right) + \tau}.$$



Effective phase shift

$$\Delta\phi \equiv \max_{\omega}\{\alpha(\omega)\} - \min_{\omega}\{\alpha(\omega)\} \implies \Delta\phi = \pi \text{ for } \tau = 1/2$$

Transmission properties



- Phase dependent symmetries

$$\phi_0 = \pi/2 \quad \mathcal{T}(\omega - \omega_0) = 1 - \mathcal{T}(-[\omega - \omega_0])$$

$$\phi_0 = 0 \quad \mathcal{T}(\omega - \omega_0) = \mathcal{T}(-[\omega - \omega_0])$$

- Filter analogy, $\tau > 0.5$, $\phi_0 = \pi/2$

$$R = 1/2 - \sqrt{\tau(1 - \tau)} \quad \text{Ripple}$$

$$\Delta\omega = \frac{\Delta}{\pi} \left[\arcsin \left(\frac{1}{\sqrt{2\tau}} \right) - \frac{\pi}{4} \right] \quad \text{Transition width, roll-off}$$

Thermoelectric scattering theory

Linear response theory (non-interacting), charge and heat currents

Butcher, 1990

$$\begin{pmatrix} I_e \\ I_h \end{pmatrix} = \begin{pmatrix} \mathcal{L}_{eV} & \mathcal{L}_{eT} \\ \mathcal{L}_{hV} & \mathcal{L}_{hT} \end{pmatrix} \begin{pmatrix} F_V \\ F_T \end{pmatrix}$$

Thermodynamic forces

$$F_V = eV/k_B T \quad F_T = \Delta T / (k_B T)^2$$

$$\text{where } V = V_H - V_C, \quad \Delta T = T_H - T_C$$

Onsager matrix

$$\begin{pmatrix} \mathcal{L}_{eV} & \mathcal{L}_{eT} \\ \mathcal{L}_{hV} & \mathcal{L}_{hT} \end{pmatrix} = \frac{1}{h} \int d\omega \mathcal{T}(\omega) \xi(\omega) \begin{pmatrix} e & e\omega \\ \omega & \omega^2 \end{pmatrix}$$

$$\text{where } \xi(\omega) = \left(2 \cosh \frac{\omega}{2} \right)^{-2}$$

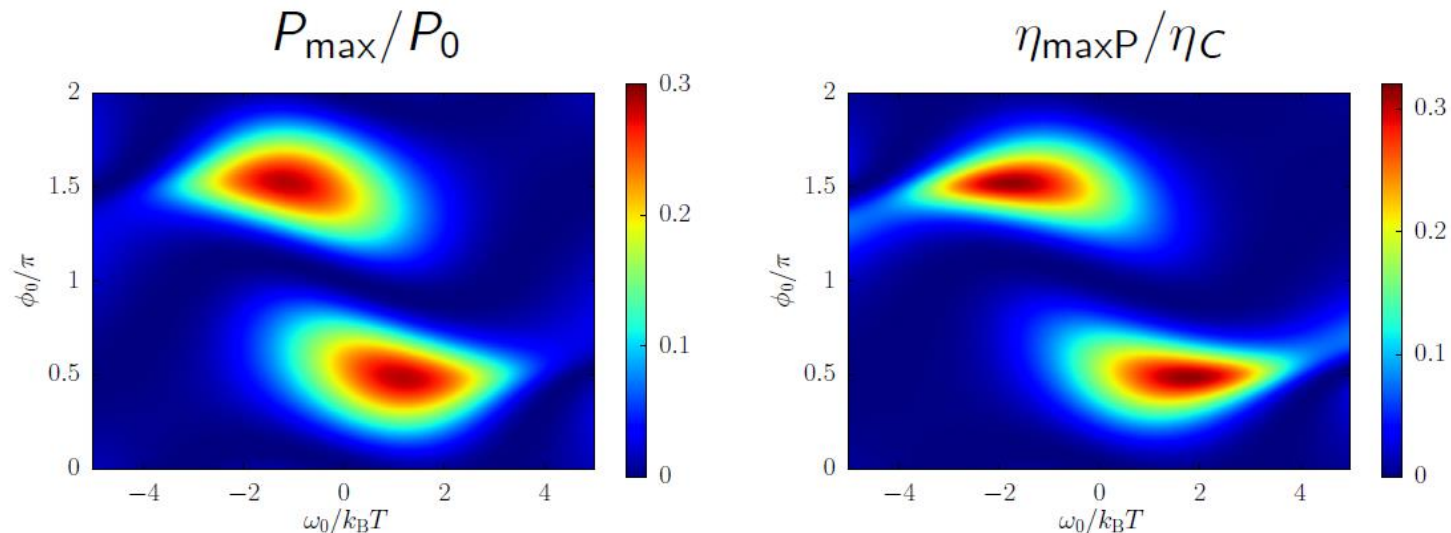
Power and efficiency

Maximum power generated, with respect to voltage

$$P_{\max} = \frac{k_B T}{4e} \frac{(\mathcal{L}_{eT})^2}{\mathcal{L}_{eV}} (F_T)^2.$$

Efficiency at maximum power

$$\eta_{\max P} = \frac{P_{\max}}{I_h} = \frac{\eta_C}{2e} \frac{(\mathcal{L}_{eT})^2}{2\mathcal{L}_{eV}\mathcal{L}_{hT} - \mathcal{L}_{eT}} \quad P_0 = \frac{(k_B \Delta T)^2}{h}$$



Close-to-optimal performance

Optimal, single mode performance (step function transmission)

$$\mathcal{T}(\omega) = \theta(\omega - \omega_0) \implies$$

$$P_{\max} = 0.32 \frac{(k_B \Delta T)^2}{h} \quad \text{for} \quad \omega_0 = 1.16 k_B T$$

$$\eta_{\max P} = 0.35 \eta_C$$

Numerical optimization

Optimizing $\mathcal{T}(\omega)$ over $0 \leq \tau \leq 1, 0 \leq \phi_0 \leq \pi, \omega_0, \Delta$

- Parameters $\tau = 0.61, \phi_0 = 0.52\pi, \omega_0 = 1.17 k_B T, \Delta = 24.2 k_B T$

- Values

$$P_{\max} = 0.285 \frac{(k_B \Delta T)^2}{h} \quad \eta_{\max P} = 0.29 \eta_C$$

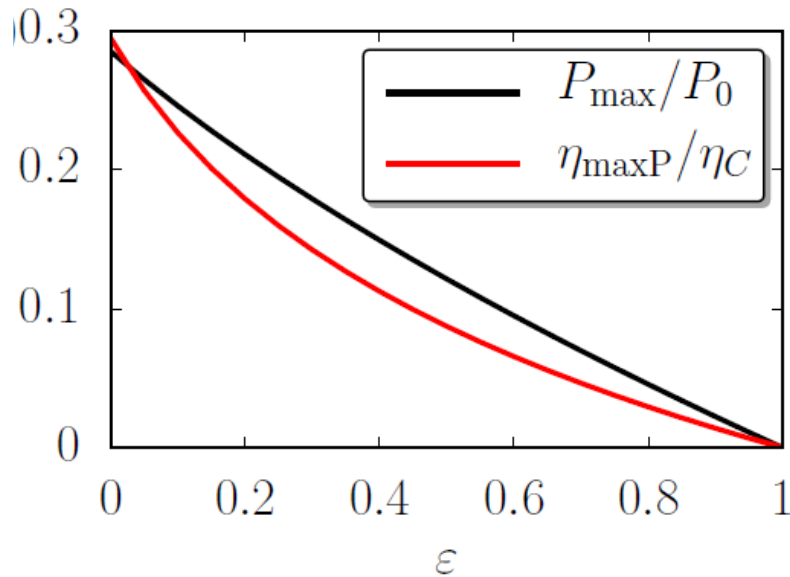
90%

← of optimal →

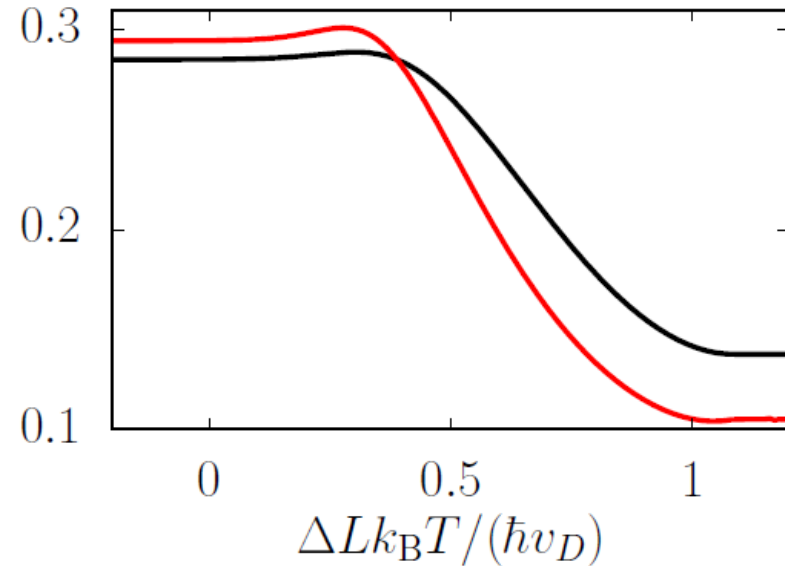
83%

Show-stoppers?

Dephasing



Asymmetry



- Survives for moderate dephasing strength $0 \leq \epsilon \leq 1$
- Zero for complete dephasing.
- Not very sensitive to arm length asymmetry $\Delta L = L_1 - L_2$

Conclusions

- Interference-only thermoelectrics potentially optimal
- Close-to-optimal performance in edge-state setup
- Not very sensitive to dephasing and asymmetry

